

Entanglement entropies in minimal models from null-vector equations

Yacine Ikhlef
LPTHE, Université Paris-6/CNRS

October 2017
Lyon
in honour of Jean-Michel Maillet

This talk is based on:

Thomas Dupic, Benoît Estienne and YI

Entanglement entropies in minimal models from null vectors

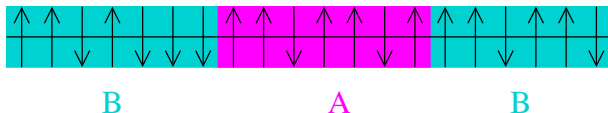
arXiv:1709.09270

Outline

1. Introduction
2. The null-vector approach
3. Entanglement entropies in the Yang-Lee model

1. Introduction

Entanglement entropies in quantum systems

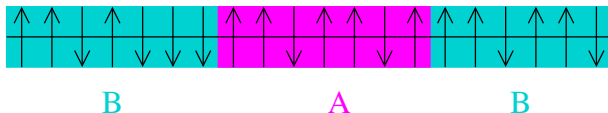


► Density matrix of the whole system $A + B$

► Pure state $|\psi\rangle \Rightarrow \rho = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle}$

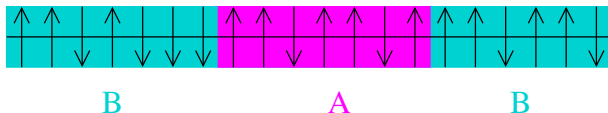
► Mixed state at temperature $\beta \Rightarrow \rho = \frac{1}{Z} \exp(-\beta H)$

Entanglement entropies in quantum systems



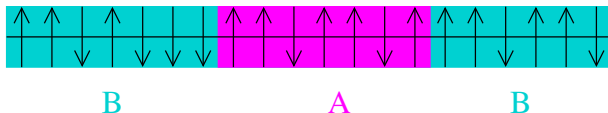
- ▶ Density matrix of the whole system $A + B$
 - ▶ Pure state $|\psi\rangle \Rightarrow \rho = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle}$
 - ▶ Mixed state at temperature $\beta \Rightarrow \rho = \frac{1}{Z} \exp(-\beta H)$
- ▶ Reduced density matrix of subsystem A : $\rho_A = \text{Tr}_B(\rho)$

Entanglement entropies in quantum systems



- ▶ Density matrix of the whole system $A + B$
 - ▶ Pure state $|\psi\rangle \Rightarrow \rho = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle}$
 - ▶ Mixed state at temperature $\beta \Rightarrow \rho = \frac{1}{Z} \exp(-\beta H)$
- ▶ Reduced density matrix of subsystem A : $\rho_A = \text{Tr}_B(\rho)$
- ▶ Von Neumann entropy: $S(A) = -\text{Tr}_A(\rho_A \log \rho_A)$

Entanglement entropies in quantum systems



- ▶ Density matrix of the whole system $A + B$
 - ▶ Pure state $|\psi\rangle \Rightarrow \rho = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle}$
 - ▶ Mixed state at temperature $\beta \Rightarrow \rho = \frac{1}{Z} \exp(-\beta H)$
- ▶ Reduced density matrix of subsystem A : $\rho_A = \text{Tr}_B(\rho)$
- ▶ Von Neumann entropy: $S(A) = -\text{Tr}_A(\rho_A \log \rho_A)$
- ▶ Rényi entropy: $S_N(A) = \frac{1}{1-N} \log \text{Tr}_A(\rho_A^N)$

Entanglement entropies in quantum systems

- ▶ Example: two $1/2$ -spins A, B

Entanglement entropies in quantum systems

- ▶ Example: two 1/2-spins A, B

- ▶ Product state: $|\psi\rangle = |\psi_A\rangle \otimes |\psi_B\rangle$

$$\rho_A = |\psi_A\rangle\langle\psi_A| \quad \Rightarrow \quad S(A) = 0$$

Region A is effectively in a pure state $|\psi_A\rangle$.

Entanglement entropies in quantum systems

- ▶ Example: two 1/2-spins A, B

- ▶ Product state: $|\psi\rangle = |\psi_A\rangle \otimes |\psi_B\rangle$

- $\rho_A = |\psi_A\rangle\langle\psi_A| \Rightarrow S(A) = 0$

- Region A is effectively in a pure state $|\psi_A\rangle$.

- ▶ Entangled state $\psi = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + e^{i\phi}|\downarrow\uparrow\rangle)$

- $\rho_A = \frac{1}{2}(|\uparrow\rangle\langle\uparrow| + |\downarrow\rangle\langle\downarrow|) \Rightarrow S(A) = \log 2$

- Region A is effectively in a thermal state.

Entanglement entropies in quantum systems

- ▶ Example: two 1/2-spins A, B

- ▶ Product state: $|\psi\rangle = |\psi_A\rangle \otimes |\psi_B\rangle$

$$\rho_A = |\psi_A\rangle\langle\psi_A| \Rightarrow S(A) = 0$$

Region A is effectively in a pure state $|\psi_A\rangle$.

- ▶ Entangled state $\psi = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + e^{i\phi}|\downarrow\uparrow\rangle)$

$$\rho_A = \frac{1}{2}(|\uparrow\rangle\langle\uparrow| + |\downarrow\rangle\langle\downarrow|) \Rightarrow S(A) = \log 2$$

Region A is effectively in a thermal state.

- ▶ Area law:

For $d \geq 1 + 1$, the entropy in the groundstate of a gapped, short-range Hamiltonian scales generally as

$$S(A) \propto \text{Area}(\partial A).$$

Entanglement entropies in quantum systems

- ▶ Example: two 1/2-spins A, B

- ▶ Product state: $|\psi\rangle = |\psi_A\rangle \otimes |\psi_B\rangle$

$$\rho_A = |\psi_A\rangle\langle\psi_A| \Rightarrow S(A) = 0$$

Region A is effectively in a pure state $|\psi_A\rangle$.

- ▶ Entangled state $\psi = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + e^{i\phi}|\downarrow\uparrow\rangle)$

$$\rho_A = \frac{1}{2}(|\uparrow\rangle\langle\uparrow| + |\downarrow\rangle\langle\downarrow|) \Rightarrow S(A) = \log 2$$

Region A is effectively in a thermal state.

- ▶ Area law:

For $d \geq 1 + 1$, the entropy in the groundstate of a gapped, short-range Hamiltonian scales generally as

$$S(A) \propto \text{Area}(\partial A).$$

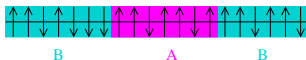
- ▶ \neq Gapless systems in $d = 1 + 1$: **Conformal Field Theory**
[Holzhey-Larsen-Wilczek '94, Calabrese-Cardy '04]

$$S(A) \sim \frac{c}{3} \log \ell_A$$

The path-integral formalism for Rényi entropies

[Calabrese-Cardy '04]

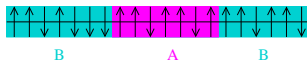
► $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$:



The path-integral formalism for Rényi entropies

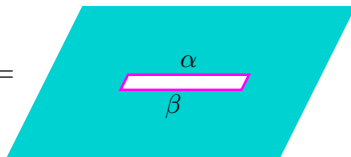
[Calabrese-Cardy '04]

► $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$:



► Reduced density matrix

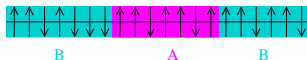
$$(\rho_A)_{\alpha\beta} =$$



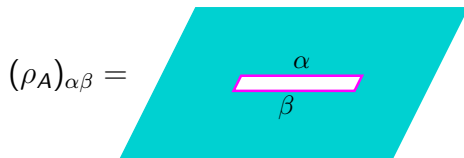
The path-integral formalism for Rényi entropies

[Calabrese-Cardy '04]

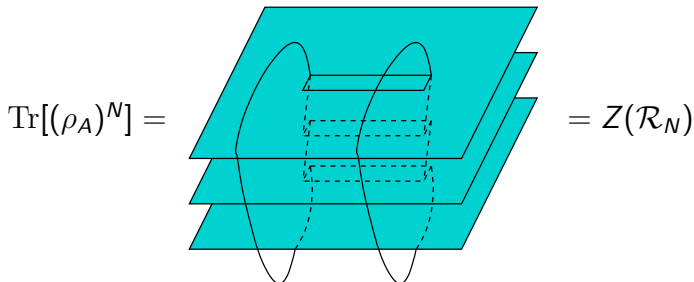
► $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$:



► Reduced density matrix



► Rényi entropy



Partition functions on Riemann surfaces

- ▶ Rényi entropy: $S_N = \frac{1}{1-N} \log \frac{Z(\mathcal{R}_N)}{Z^N}$

Partition functions on Riemann surfaces

► Rényi entropy: $S_N = \frac{1}{1-N} \log \frac{Z(\mathcal{R}_N)}{Z^N}$

► One interval $[u, v]$:
$$\begin{cases} \mathcal{R}_N & \rightarrow \mathbb{C} \cup \{\infty\} \\ z & \mapsto w = \left(\frac{z-u}{z-v} \right)^{1/N} \end{cases}$$

Partition functions on Riemann surfaces

► Rényi entropy: $S_N = \frac{1}{1-N} \log \frac{Z(\mathcal{R}_N)}{Z^N}$

► One interval $[u, v]$:
$$\begin{cases} \mathcal{R}_N & \rightarrow \mathbb{C} \cup \{\infty\} \\ z & \mapsto w = \left(\frac{z-u}{z-v} \right)^{1/N} \end{cases}$$

► Generic case:

$$\begin{cases} A = \text{union of } p \text{ intervals} \\ N \text{ copies} \end{cases} \Rightarrow \text{genus}(\mathcal{R}_N) = (N-1)(p-1)$$

Partition functions on Riemann surfaces

► Rényi entropy: $S_N = \frac{1}{1-N} \log \frac{Z(\mathcal{R}_N)}{Z^N}$

► One interval $[u, v]$:
$$\begin{cases} \mathcal{R}_N & \rightarrow \mathbb{C} \cup \{\infty\} \\ z & \mapsto w = \left(\frac{z-u}{z-v} \right)^{1/N} \end{cases}$$

► Generic case:

$$\begin{cases} A = \text{union of } p \text{ intervals} \\ N \text{ copies} \end{cases} \Rightarrow \text{genus}(\mathcal{R}_N) = (N-1)(p-1)$$

► Example: $p = 2, N = 4$



Partition functions on Riemann surfaces

► Rényi entropy: $S_N = \frac{1}{1-N} \log \frac{Z(\mathcal{R}_N)}{Z^N}$

► One interval $[u, v]$:
$$\begin{cases} \mathcal{R}_N & \rightarrow \mathbb{C} \cup \{\infty\} \\ z & \mapsto w = \left(\frac{z-u}{z-v} \right)^{1/N} \end{cases}$$

► Generic case:

$$\begin{cases} A = \text{union of } p \text{ intervals} \\ N \text{ copies} \end{cases} \Rightarrow \text{genus}(\mathcal{R}_N) = (N-1)(p-1)$$

► Example: $p = 2, N = 4$



► For $c = 1$: results available from [Zamolodchikov '87], [Dixon, Friedan, Martinec, Shenker '87], [Alvarez-Gaumé, Gost, Moore, Nelson, Vafa '87], [Dijkgraaf, Verlinde, Verlinde '88]

Overview of Entanglement Entropies in CFT

- ▶ Scaling argument $S \propto \frac{c}{3} \log \ell$ [Holzhey, Larsen, Wilczek '94]:
- ▶ Path-integral approach, EE by conformal mapping for $A = [u, v]$ [Calabrese, Cardy '04]
- ▶ Compute EE at $g > 0$ for $c = 1$ and/or Ising CFTs [Calabrese, Cardy, Tonni, Tagliacozzo, Alba, Misguich, Pasquier, Stéphan, Furukawa, Shiraishi, Essler, Campostrini, Nienhuis '06–'12]
- ▶ EE for excited states [Sierra, Alcaraz, Berganza, Palmai '12–'16]
- ▶ EE for integrable QFTs [Castro-Alvaredo, Doyon, Cardy, Blondeau-Fournier '07–'15]:
- ▶ EE, entanglement spectrum, fidelity using CTM [Franchini, Its, Korepin, Takhtajan, Evangelisti, Weston '11–'12]
- ▶ Entanglement after a quench [Cardy '11]
- ▶ EE for non-unitary CFTs [Castro-Alvaredo, Doyon, Ravanini, Bianchini, Levi, Couvreur, Jacobsen, Saleur '14–'17]
- ▶ Entanglement spectra in FQHE [Li, Haldane, Read, Rezayi, Dubail, Eisler, Peschel, Cardy, Tonni '08–'17]:

▶ ...

2. The null-vector approach

The \mathbb{Z}_N cyclic orbifold CFT

[Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ “Replicate the degrees of freedom instead of the surface”

The \mathbb{Z}_N cyclic orbifold CFT

[Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ “Replicate the degrees of freedom instead of the surface”
- ▶ Hilbert space: $\mathcal{H}_{\text{orb}} = (\mathcal{H}_{\text{CFT}})^{\otimes N} / \mathbb{Z}_N$

The \mathbb{Z}_N cyclic orbifold CFT

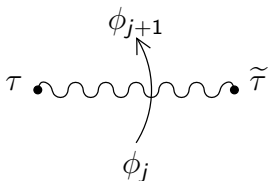
[Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ “Replicate the degrees of freedom instead of the surface”
- ▶ Hilbert space: $\mathcal{H}_{\text{orb}} = (\mathcal{H}_{\text{CFT}})^{\otimes N} / \mathbb{Z}_N$
- ▶ Local configuration of fields: (ϕ_1, \dots, ϕ_N)

The \mathbb{Z}_N cyclic orbifold CFT

[Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ “Replicate the degrees of freedom instead of the surface”
- ▶ Hilbert space: $\mathcal{H}_{\text{orb}} = (\mathcal{H}_{\text{CFT}})^{\otimes N} / \mathbb{Z}_N$
- ▶ Local configuration of fields: (ϕ_1, \dots, ϕ_N)
- ▶ Twist operator inserting a branch point:

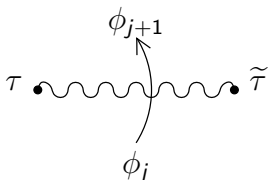


$$\tau(0) \cdot (\phi_1, \dots, \phi_{N-1}, \phi_N)(e^{2i\pi} z) = \tau(0) \cdot (\phi_2, \dots, \phi_N, \phi_1)(z)$$

The \mathbb{Z}_N cyclic orbifold CFT

[Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ “Replicate the degrees of freedom instead of the surface”
- ▶ Hilbert space: $\mathcal{H}_{\text{orb}} = (\mathcal{H}_{\text{CFT}})^{\otimes N} / \mathbb{Z}_N$
- ▶ Local configuration of fields: (ϕ_1, \dots, ϕ_N)
- ▶ Twist operator inserting a branch point:



$$\tau(0) \cdot (\phi_1, \dots, \phi_{N-1}, \phi_N)(e^{2i\pi} z) = \tau(0) \cdot (\phi_2, \dots, \phi_N, \phi_1)(z)$$

- ▶ Examples:

- ▶ $\langle \tau(u_1) \tilde{\tau}(v_1) \dots \tau(u_p) \tilde{\tau}(v_p) \rangle$
- ▶ $\langle \Phi(\infty) \tau(u) \tilde{\tau}(v) \Phi(0) \rangle, \quad \Phi := \phi_{12} \otimes \dots \otimes \phi_{12}$

The \mathbb{Z}_N cyclic orbifold CFT

Generating algebra [Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ Family of currents $\widehat{T}^{(r)}(z) = \sum_{j=1}^N e^{2i\pi rj/N} T_j(z)$

The \mathbb{Z}_N cyclic orbifold CFT

Generating algebra [Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ Family of currents $\hat{T}^{(r)}(z) = \sum_{j=1}^N e^{2i\pi rj/N} T_j(z)$
- ▶ $\hat{T}^{(0)}(z)$ =energy-momentum, $\hat{T}^{(r \neq 0)}(z)$ =additional currents

The \mathbb{Z}_N cyclic orbifold CFT

Generating algebra [Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ Family of currents $\widehat{T}^{(r)}(z) = \sum_{j=1}^N e^{2i\pi rj/N} T_j(z)$
- ▶ $\widehat{T}^{(0)}(z)$ =energy-momentum, $\widehat{T}^{(r \neq 0)}(z)$ =additional currents
- ▶ Simple monodromy: $\tau(0) \cdot \widehat{T}^{(r)}(e^{2i\pi} z) = e^{-\frac{2i\pi r}{N}} \times \tau(0) \cdot \widehat{T}^{(r)}(z)$

The \mathbb{Z}_N cyclic orbifold CFT

Generating algebra [Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ Family of currents $\hat{T}^{(r)}(z) = \sum_{j=1}^N e^{2i\pi rj/N} T_j(z)$
- ▶ $\hat{T}^{(0)}(z)$ =energy-momentum, $\hat{T}^{(r \neq 0)}(z)$ =additional currents
- ▶ Simple monodromy: $\tau(0) \cdot \hat{T}^{(r)}(e^{2i\pi} z) = e^{-\frac{2i\pi r}{N}} \times \tau(0) \cdot \hat{T}^{(r)}(z)$
- ▶ Fourier modes: $\hat{L}_m^{(r)} = \frac{1}{2i\pi} \oint dz z^{m+1} \hat{T}^{(r)}(z)$ $m \in \mathbb{Z}/N$

The \mathbb{Z}_N cyclic orbifold CFT

Generating algebra [Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ Family of currents $\widehat{T}^{(r)}(z) = \sum_{j=1}^N e^{2i\pi rj/N} T_j(z)$
- ▶ $\widehat{T}^{(0)}(z)$ =energy-momentum, $\widehat{T}^{(r \neq 0)}(z)$ =additional currents
- ▶ Simple monodromy: $\tau(0) \cdot \widehat{T}^{(r)}(e^{2i\pi} z) = e^{-\frac{2i\pi r}{N}} \times \tau(0) \cdot \widehat{T}^{(r)}(z)$
- ▶ Fourier modes: $\widehat{L}_m^{(r)} = \frac{1}{2i\pi} \oint dz z^{m+1} \widehat{T}^{(r)}(z)$ $m \in \mathbb{Z}/N$
- ▶ Orbifold Virasoro algebra:

$$[\widehat{L}_m^{(r)}, \widehat{L}_n^{(s)}] = (m - n) \widehat{L}_{m+n}^{(r+s)} + \frac{Nc}{12} m(m^2 - 1) \delta_{m+n,0} \delta_{r+s,0}$$

The \mathbb{Z}_N cyclic orbifold CFT

Generating algebra [Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ Family of currents $\widehat{T}^{(r)}(z) = \sum_{j=1}^N e^{2i\pi rj/N} T_j(z)$
- ▶ $\widehat{T}^{(0)}(z)$ =energy-momentum, $\widehat{T}^{(r \neq 0)}(z)$ =additional currents
- ▶ Simple monodromy: $\tau(0) \cdot \widehat{T}^{(r)}(e^{2i\pi} z) = e^{-\frac{2i\pi r}{N}} \times \tau(0) \cdot \widehat{T}^{(r)}(z)$
- ▶ Fourier modes: $\widehat{L}_m^{(r)} = \frac{1}{2i\pi} \oint dz z^{m+1} \widehat{T}^{(r)}(z)$ $m \in \mathbb{Z}/N$
- ▶ Orbifold Virasoro algebra:
$$\left[\widehat{L}_m^{(r)}, \widehat{L}_n^{(s)} \right] = (m - n) \widehat{L}_{m+n}^{(r+s)} + \frac{Nc}{12} m(m^2 - 1) \delta_{m+n,0} \delta_{r+s,0}$$
- ▶ Operator content

The \mathbb{Z}_N cyclic orbifold CFT

Generating algebra [Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ Family of currents $\widehat{T}^{(r)}(z) = \sum_{j=1}^N e^{2i\pi rj/N} T_j(z)$
- ▶ $\widehat{T}^{(0)}(z)$ =energy-momentum, $\widehat{T}^{(r \neq 0)}(z)$ =additional currents
- ▶ Simple monodromy: $\tau(0) \cdot \widehat{T}^{(r)}(e^{2i\pi} z) = e^{-\frac{2i\pi r}{N}} \times \tau(0) \cdot \widehat{T}^{(r)}(z)$
- ▶ Fourier modes: $\widehat{L}_m^{(r)} = \frac{1}{2i\pi} \oint dz z^{m+1} \widehat{T}^{(r)}(z)$ $m \in \mathbb{Z}/N$
- ▶ Orbifold Virasoro algebra:
$$\left[\widehat{L}_m^{(r)}, \widehat{L}_n^{(s)} \right] = (m - n) \widehat{L}_{m+n}^{(r+s)} + \frac{Nc}{12} m(m^2 - 1) \delta_{m+n,0} \delta_{r+s,0}$$
- ▶ Operator content
 - ▶ Untwisted primary fields: $\phi_1 \otimes \cdots \otimes \phi_N$, $h = h_1 + \cdots + h_N$

The \mathbb{Z}_N cyclic orbifold CFT

Generating algebra [Crnković-Sotkov-Stanishkov '89, Klemm-Schmidt '90, Borisov-Halpern-Schweigert '98]

- ▶ Family of currents $\widehat{T}^{(r)}(z) = \sum_{j=1}^N e^{2i\pi rj/N} T_j(z)$
- ▶ $\widehat{T}^{(0)}(z)$ =energy-momentum, $\widehat{T}^{(r \neq 0)}(z)$ =additional currents
- ▶ Simple monodromy: $\tau(0) \cdot \widehat{T}^{(r)}(e^{2i\pi} z) = e^{-\frac{2i\pi r}{N}} \times \tau(0) \cdot \widehat{T}^{(r)}(z)$

- ▶ Fourier modes: $\widehat{L}_m^{(r)} = \frac{1}{2i\pi} \oint dz z^{m+1} \widehat{T}^{(r)}(z)$ $m \in \mathbb{Z}/N$

- ▶ Orbifold Virasoro algebra:

$$\left[\widehat{L}_m^{(r)}, \widehat{L}_n^{(s)} \right] = (m - n) \widehat{L}_{m+n}^{(r+s)} + \frac{Nc}{12} m(m^2 - 1) \delta_{m+n,0} \delta_{r+s,0}$$

- ▶ Operator content

- ▶ Untwisted primary fields: $\phi_1 \otimes \cdots \otimes \phi_N$, $h = h_1 + \cdots + h_N$
- ▶ Twisted primary fields: $\tau_\phi = : \tau \phi :$

Induction procedure

[Crnković-Sotkov-Stanishkov '89, Borisov-Halpern-Schweigert '98]

- ▶ Consider (radial) quantisation around a branch point

Induction procedure

[Crnković-Sotkov-Stanishkov '89, Borisov-Halpern-Schweigert '98]

- ▶ Consider (radial) quantisation around a branch point
- ▶ Conformal mapping $z \mapsto w = z^{1/N}$ in the vicinity of $z = 0$:

Induction procedure

[Crnković-Sotkov-Stanishkov '89, Borisov-Halpern-Schweigert '98]

- ▶ Consider (radial) quantisation around a branch point
- ▶ Conformal mapping $z \mapsto w = z^{1/N}$ in the vicinity of $z = 0$:

$$\text{▶ } \widehat{T}^{(r)}(z) \mapsto w^{2-2N} \sum_{j=1}^N e^{\frac{2i\pi(j+2)}{N}} T(e^{\frac{2i\pi j}{N}} w) + \frac{c}{24z^2} \left(N - \frac{1}{N}\right) \delta_{r,0}$$

Induction procedure

[Crnković-Sotkov-Stanishkov '89, Borisov-Halpern-Schweigert '98]

- ▶ Consider (radial) quantisation around a branch point
- ▶ Conformal mapping $z \mapsto w = z^{1/N}$ in the vicinity of $z = 0$:

- ▶ $\widehat{T}^{(r)}(z) \mapsto w^{2-2N} \sum_{j=1}^N e^{\frac{2i\pi(j+2)}{N}} T(e^{\frac{2i\pi j}{N}} w) + \frac{c}{24z^2} \left(N - \frac{1}{N}\right) \delta_{r,0}$

- ▶ $\widehat{L}_m^{(r)} \mapsto \frac{1}{N} L_{Nm} + \frac{c}{24} \left(N - \frac{1}{N}\right) \delta_{r,0} \delta_{m,0}$

Induction procedure

[Crnković-Sotkov-Stanishkov '89, Borisov-Halpern-Schweigert '98]

- ▶ Consider (radial) quantisation around a branch point
- ▶ Conformal mapping $z \mapsto w = z^{1/N}$ in the vicinity of $z = 0$:

$$\text{▶ } \widehat{T}^{(r)}(z) \mapsto w^{2-2N} \sum_{j=1}^N e^{\frac{2i\pi(j+2)}{N}} T(e^{\frac{2i\pi j}{N}} w) + \frac{c}{24z^2} \left(N - \frac{1}{N}\right) \delta_{r,0}$$

$$\text{▶ } \widehat{L}_m^{(r)} \mapsto \frac{1}{N} L_{Nm} + \frac{c}{24} \left(N - \frac{1}{N}\right) \delta_{r,0} \delta_{m,0}$$

$$\text{▶ } \tau_\phi(z) \mapsto w^{(1-N)h_\phi} \phi(w)$$

Induction procedure

[Crnković-Sotkov-Stanishkov '89, Borisov-Halpern-Schweigert '98]

- ▶ Consider (radial) quantisation around a branch point
- ▶ Conformal mapping $z \mapsto w = z^{1/N}$ in the vicinity of $z = 0$:
 - ▶ $\widehat{T}^{(r)}(z) \mapsto w^{2-2N} \sum_{j=1}^N e^{\frac{2i\pi(j+2)}{N}} T(e^{\frac{2i\pi j}{N}} w) + \frac{c}{24z^2} \left(N - \frac{1}{N}\right) \delta_{r,0}$
 - ▶ $\widehat{L}_m^{(r)} \mapsto \frac{1}{N} L_{Nm} + \frac{c}{24} \left(N - \frac{1}{N}\right) \delta_{r,0} \delta_{m,0}$
 - ▶ $\tau_\phi(z) \mapsto w^{(1-N)h_\phi} \phi(w)$
- ▶ Applications:

Induction procedure

[Crnković-Sotkov-Stanishkov '89, Borisov-Halpern-Schweigert '98]

- ▶ Consider (radial) quantisation around a branch point
- ▶ Conformal mapping $z \mapsto w = z^{1/N}$ in the vicinity of $z = 0$:
 - ▶ $\widehat{T}^{(r)}(z) \mapsto w^{2-2N} \sum_{j=1}^N e^{\frac{2i\pi(j+2)}{N}} T(e^{\frac{2i\pi j}{N}} w) + \frac{c}{24z^2} \left(N - \frac{1}{N}\right) \delta_{r,0}$
 - ▶ $\widehat{L}_m^{(r)} \mapsto \frac{1}{N} L_{Nm} + \frac{c}{24} \left(N - \frac{1}{N}\right) \delta_{r,0} \delta_{m,0}$
 - ▶ $\tau_\phi(z) \mapsto w^{(1-N)h_\phi} \phi(w)$
- ▶ Applications:
 1. τ_ϕ has dimension $\widehat{h}_\phi = \frac{c}{24} \left(N - \frac{1}{N}\right) + \frac{h}{N}$

Induction procedure

[Crnković-Sotkov-Stanishkov '89, Borisov-Halpern-Schweigert '98]

- ▶ Consider (radial) quantisation around a branch point
- ▶ Conformal mapping $z \mapsto w = z^{1/N}$ in the vicinity of $z = 0$:
 - ▶ $\widehat{T}^{(r)}(z) \mapsto w^{2-2N} \sum_{j=1}^N e^{\frac{2i\pi(j+2)}{N}} T(e^{\frac{2i\pi j}{N}} w) + \frac{c}{24z^2} \left(N - \frac{1}{N}\right) \delta_{r,0}$
 - ▶ $\widehat{L}_m^{(r)} \mapsto \frac{1}{N} L_{Nm} + \frac{c}{24} \left(N - \frac{1}{N}\right) \delta_{r,0} \delta_{m,0}$
 - ▶ $\tau_\phi(z) \mapsto w^{(1-N)h_\phi} \phi(w)$
- ▶ Applications:
 1. τ_ϕ has dimension $\widehat{h}_\phi = \frac{c}{24} \left(N - \frac{1}{N}\right) + \frac{h}{N}$
 2. **If ϕ degenerate at level ℓ then τ_ϕ degenerate at level ℓ/N**

Induction procedure

[Crnković-Sotkov-Stanishkov '89, Borisov-Halpern-Schweigert '98]

- ▶ Consider (radial) quantisation around a branch point
- ▶ Conformal mapping $z \mapsto w = z^{1/N}$ in the vicinity of $z = 0$:
 - ▶ $\widehat{T}^{(r)}(z) \mapsto w^{2-2N} \sum_{j=1}^N e^{\frac{2i\pi(j+2)}{N}} T(e^{\frac{2i\pi j}{N}} w) + \frac{c}{24z^2} \left(N - \frac{1}{N}\right) \delta_{r,0}$
 - ▶ $\widehat{L}_m^{(r)} \mapsto \frac{1}{N} L_{Nm} + \frac{c}{24} \left(N - \frac{1}{N}\right) \delta_{r,0} \delta_{m,0}$
 - ▶ $\tau_\phi(z) \mapsto w^{(1-N)h_\phi} \phi(w)$
- ▶ Applications:
 1. τ_ϕ has dimension $\widehat{h}_\phi = \frac{c}{24} \left(N - \frac{1}{N}\right) + \frac{h}{N}$
 2. **If ϕ degenerate at level ℓ then τ_ϕ degenerate at level ℓ/N**
- ▶ Basic example: $L_{-1}\mathbf{1} = 0 \Rightarrow \widehat{L}_{-1/N}^{(-1)}\tau = 0$

The \mathbb{Z}_N orbifold of a minimal model

- ▶ EE in minimal model $\mathcal{M}(p, q)$ with $c = 1 - \frac{6(p-q)^2}{pq}$

The \mathbb{Z}_N orbifold of a minimal model

- ▶ EE in minimal model $\mathcal{M}(p, q)$ with $c = 1 - \frac{6(p-q)^2}{pq}$
- ▶ Degenerate operators : Kac table $h_{rs} = \frac{(pr-qs)^2 - (p-q)^2}{4pq}$

The \mathbb{Z}_N orbifold of a minimal model

- ▶ EE in minimal model $\mathcal{M}(p, q)$ with $c = 1 - \frac{6(p-q)^2}{pq}$
- ▶ Degenerate operators : Kac table $h_{rs} = \frac{(pr-qs)^2 - (p-q)^2}{4pq}$
- ▶ Symmetry $\phi_{rs} \equiv \phi_{q-r, p-s} \Rightarrow$ two null vectors for ϕ_{rs}

The \mathbb{Z}_N orbifold of a minimal model

- ▶ EE in minimal model $\mathcal{M}(p, q)$ with $c = 1 - \frac{6(p-q)^2}{pq}$
- ▶ Degenerate operators : Kac table $h_{rs} = \frac{(pr-qs)^2 - (p-q)^2}{4pq}$
- ▶ Symmetry $\phi_{rs} \equiv \phi_{q-r, p-s} \Rightarrow$ two null vectors for ϕ_{rs}
- ▶ Translate into two null vectors for $\tau_{\phi_{rs}}$

The \mathbb{Z}_N orbifold of a minimal model

- ▶ EE in minimal model $\mathcal{M}(p, q)$ with $c = 1 - \frac{6(p-q)^2}{pq}$
- ▶ Degenerate operators : Kac table $h_{rs} = \frac{(pr-qs)^2 - (p-q)^2}{4pq}$
- ▶ Symmetry $\phi_{rs} \equiv \phi_{q-r, p-s} \Rightarrow$ two null vectors for ϕ_{rs}
- ▶ Translate into two null vectors for $\tau_{\phi_{rs}}$

- ▶ Example: Ising = $\mathcal{M}(3, 4) =$

| | | |
|------------|----------|------------|
| ϵ | σ | 1 |
| 1 | σ | ϵ |

The \mathbb{Z}_N orbifold of a minimal model

- ▶ EE in minimal model $\mathcal{M}(p, q)$ with $c = 1 - \frac{6(p-q)^2}{pq}$
- ▶ Degenerate operators : Kac table $h_{rs} = \frac{(pr-qs)^2 - (p-q)^2}{4pq}$
- ▶ Symmetry $\phi_{rs} \equiv \phi_{q-r, p-s} \Rightarrow$ two null vectors for ϕ_{rs}
- ▶ Translate into two null vectors for $\tau_{\phi_{rs}}$

- ▶ Example: Ising = $\mathcal{M}(3, 4) =$

| | | |
|--------------|----------|--------------|
| ϵ | σ | $\mathbf{1}$ |
| $\mathbf{1}$ | σ | ϵ |

- ▶
$$\begin{cases} L_{-1}\mathbf{1} = 0 \\ (L_{-6} + aL_{-4}L_{-2} + bL_{-3}^2 + cL_{-2}^3)\mathbf{1} = 0 \end{cases}$$

The \mathbb{Z}_N orbifold of a minimal model

- ▶ EE in minimal model $\mathcal{M}(p, q)$ with $c = 1 - \frac{6(p-q)^2}{pq}$
- ▶ Degenerate operators : Kac table $h_{rs} = \frac{(pr-qs)^2 - (p-q)^2}{4pq}$
- ▶ Symmetry $\phi_{rs} \equiv \phi_{q-r, p-s} \Rightarrow$ two null vectors for ϕ_{rs}
- ▶ Translate into two null vectors for $\tau_{\phi_{rs}}$

- ▶ Example: Ising = $\mathcal{M}(3, 4) =$

| | | |
|--------------|----------|--------------|
| ϵ | σ | $\mathbf{1}$ |
| $\mathbf{1}$ | σ | ϵ |

$$\begin{cases} L_{-1}\mathbf{1} = 0 \\ (L_{-6} + aL_{-4}L_{-2} + bL_{-3}^2 + cL_{-2}^3)\mathbf{1} = 0 \end{cases}$$

$$\Rightarrow \begin{cases} \widehat{L}_{-1/N} \tau = 0 \\ \left[N\widehat{L}_{-\frac{6}{N}} + aN^2\widehat{L}_{-\frac{4}{N}}\widehat{L}_{-\frac{2}{N}} + bN^2\widehat{L}_{-\frac{3}{N}}^2 + cN^3\widehat{L}_{-\frac{2}{N}}^3 \right] \tau = 0 \end{cases}$$

From null vectors to differential equations

- ▶ Four-point functions in terms of $x = \frac{(z_1 - z_2)(z_3 - z_4)}{(z_1 - z_3)(z_2 - z_4)}$:

$$\langle \mathcal{O}_1(z_1, \bar{z}_1) \dots \mathcal{O}_4(z_4, \bar{z}_4) \rangle = (\dots) \langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) \mathcal{O}_4(0) \rangle$$

From null vectors to differential equations

- ▶ Four-point functions in terms of $x = \frac{(z_1 - z_2)(z_3 - z_4)}{(z_1 - z_3)(z_2 - z_4)}$:

$$\langle \mathcal{O}_1(z_1, \bar{z}_1) \dots \mathcal{O}_4(z_4, \bar{z}_4) \rangle = (\dots) \langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) \mathcal{O}_4(0) \rangle$$

- ▶ In standard CFT, action of L_{-m} for $m \in \mathbb{N}$: [BPZ '84]

$$\langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) (L_{-m} \mathcal{O}_4)(0) \rangle = \mathcal{L}_{-m} \langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) \mathcal{O}_4(0) \rangle$$

From null vectors to differential equations

- ▶ Four-point functions in terms of $x = \frac{(z_1 - z_2)(z_3 - z_4)}{(z_1 - z_3)(z_2 - z_4)}$:

$$\langle \mathcal{O}_1(z_1, \bar{z}_1) \dots \mathcal{O}_4(z_4, \bar{z}_4) \rangle = (\dots) \langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) \mathcal{O}_4(0) \rangle$$

- ▶ In standard CFT, action of L_{-m} for $m \in \mathbb{N}$: [BPZ '84]

$$\langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) (L_{-m} \mathcal{O}_4)(0) \rangle = \mathcal{L}_{-m} \langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) \mathcal{O}_4(0) \rangle$$

- ▶ In orbifold Virasoro algebra: action of $\widehat{L}_{-m}^{(r)}$?

From null vectors to differential equations

- ▶ Four-point functions in terms of $x = \frac{(z_1 - z_2)(z_3 - z_4)}{(z_1 - z_3)(z_2 - z_4)}$:

$$\langle \mathcal{O}_1(z_1, \bar{z}_1) \dots \mathcal{O}_4(z_4, \bar{z}_4) \rangle = (\dots) \langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) \mathcal{O}_4(0) \rangle$$

- ▶ In standard CFT, action of L_{-m} for $m \in \mathbb{N}$: [BPZ '84]

$$\langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) (L_{-m} \mathcal{O}_4)(0) \rangle = \mathcal{L}_{-m} \langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) \mathcal{O}_4(0) \rangle$$

- ▶ In orbifold Virasoro algebra: action of $\widehat{L}_{-m}^{(r)}$?
- ▶ Strategy: obtain linear relation between

$$\begin{aligned} & \langle (\widehat{L}_{m_1}^{(r_1)} \mathcal{O}_1) \mathcal{O}_2 \mathcal{O}_3 \mathcal{O}_4 \rangle, \langle \mathcal{O}_1 (\widehat{L}_{m_2}^{(r_2)} \mathcal{O}_2) \mathcal{O}_3 \mathcal{O}_4 \rangle, \\ & \langle \mathcal{O}_1 \mathcal{O}_2 (\widehat{L}_{m_3}^{(r_3)} \mathcal{O}_3) \mathcal{O}_4 \rangle, \langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3 (\widehat{L}_{m_4}^{(r_4)} \mathcal{O}_4) \rangle \end{aligned}$$

using **orbifold Ward identities** = Closed-contour relations:

$$\oint_C dz (z-1)^{m_2+1} (z-x)^{m_3+1} z^{m_4+1} \times \langle \mathcal{O}_1(\infty) \mathcal{O}_2(1) \mathcal{O}_3(x, \bar{x}) \widehat{T}^{(r)}(z) \mathcal{O}_4(0) \rangle = 0$$

where $\mathcal{O}_j \in [kj] \Rightarrow m_j \in \mathbb{Z} + rk_j/N$

Summary of the null-vector approach

1. Consider eigenstate $|\psi\rangle$ of $H_{\mathcal{M}(p,q)}$
 - ▶ Write Rényi EE $S_N(A) = \frac{1}{1-N} \log G(x, \bar{x})$
 - ▶ Example: $G(x, \bar{x}) = \langle \Psi(\infty) \tau(1) \tilde{\tau}(x, \bar{x}) \Psi(0) \rangle_{\text{orb}}$

Summary of the null-vector approach

1. Consider eigenstate $|\psi\rangle$ of $H_{\mathcal{M}(p,q)}$
 - ▶ Write Rényi EE $S_N(A) = \frac{1}{1-N} \log G(x, \bar{x})$
 - ▶ Example: $G(x, \bar{x}) = \langle \Psi(\infty) \tau(1) \tilde{\tau}(x, \bar{x}) \Psi(0) \rangle_{\text{orb}}$
2. Induction: null vectors for $\mathbf{1} \Rightarrow$ null vectors for τ

Summary of the null-vector approach

1. Consider eigenstate $|\psi\rangle$ of $H_{\mathcal{M}(p,q)}$
 - ▶ Write Rényi EE $S_N(A) = \frac{1}{1-N} \log G(x, \bar{x})$
 - ▶ Example: $G(x, \bar{x}) = \langle \Psi(\infty) \tau(1) \tilde{\tau}(x, \bar{x}) \Psi(0) \rangle_{\text{orb}}$
2. Induction: null vectors for $\mathbf{1} \Rightarrow$ null vectors for τ
3. $\left\{ \begin{array}{l} \text{Null-vector conditions for } \tau \\ \text{Orbifold Ward id.} \end{array} \right. \Rightarrow$ diff. equation for $G(x, \bar{x})$

Summary of the null-vector approach

1. Consider eigenstate $|\psi\rangle$ of $H_{\mathcal{M}(p,q)}$
 - ▶ Write Rényi EE $S_N(A) = \frac{1}{1-N} \log G(x, \bar{x})$
 - ▶ Example: $G(x, \bar{x}) = \langle \Psi(\infty) \tau(1) \tilde{\tau}(x, \bar{x}) \Psi(0) \rangle_{\text{orb}}$
2. Induction: null vectors for $\mathbf{1} \Rightarrow$ null vectors for τ
3. $\left\{ \begin{array}{l} \text{Null-vector conditions for } \tau \\ \text{Orbifold Ward id.} \end{array} \right. \Rightarrow$ diff. equation for $G(x, \bar{x})$
4. Construct power series for holomorphic solutions:
 $\{I_j(x)\}$ for $|x| < 1$, $\{J_j(x)\}$ for $|1-x| < 1$

Summary of the null-vector approach

1. Consider eigenstate $|\psi\rangle$ of $H_{\mathcal{M}(p,q)}$
 - ▶ Write Rényi EE $S_N(A) = \frac{1}{1-N} \log G(x, \bar{x})$
 - ▶ Example: $G(x, \bar{x}) = \langle \Psi(\infty) \tau(1) \tilde{\tau}(x, \bar{x}) \Psi(0) \rangle_{\text{orb}}$
2. Induction: null vectors for $\mathbf{1} \Rightarrow$ null vectors for τ
3. $\left\{ \begin{array}{l} \text{Null-vector conditions for } \tau \\ \text{Orbifold Ward id.} \end{array} \right. \Rightarrow$ diff. equation for $G(x, \bar{x})$
4. Construct power series for holomorphic solutions:
 $\{I_i(x)\}$ for $|x| < 1$, $\{J_j(x)\}$ for $|1-x| < 1$
5. Assume conformal-block decomposition

$$G(x, \bar{x}) = \sum_{i=1}^M X_i |I_i(x)|^2 = \sum_{j=1}^M Y_j |J_j(x)|^2$$

Summary of the null-vector approach

1. Consider eigenstate $|\psi\rangle$ of $H_{\mathcal{M}(p,q)}$
 - ▶ Write Rényi EE $S_N(A) = \frac{1}{1-N} \log G(x, \bar{x})$
 - ▶ Example: $G(x, \bar{x}) = \langle \Psi(\infty) \tau(1) \tilde{\tau}(x, \bar{x}) \Psi(0) \rangle_{\text{orb}}$
2. Induction: null vectors for $\mathbf{1} \Rightarrow$ null vectors for τ
3. $\left\{ \begin{array}{l} \text{Null-vector conditions for } \tau \\ \text{Orbifold Ward id.} \end{array} \right. \Rightarrow$ diff. equation for $G(x, \bar{x})$
4. Construct power series for holomorphic solutions:
 $\{I_i(x)\}$ for $|x| < 1$, $\{J_j(x)\}$ for $|1-x| < 1$
5. Assume conformal-block decomposition

$$G(x, \bar{x}) = \sum_{i=1}^M X_i |I_i(x)|^2 = \sum_{j=1}^M Y_j |J_j(x)|^2$$

6. Solve consistency for $\{X_i\} \leftrightarrow \{Y_j\}$ (“conformal bootstrap”)

3. Entanglement entropies in the Yang-Lee model

The Yang-Lee edge singularity

[Yang-Lee '52]

- ▶ Classical Ising model in magnetic field :

$$Z(J, H) = \sum_{\{\sigma\}} \exp \left(J \sum_{\langle i, j \rangle} \sigma_i \sigma_j + H \sum_j \sigma_j \right)$$

The Yang-Lee edge singularity

[Yang-Lee '52]

- ▶ Classical Ising model in magnetic field :

$$Z(J, H) = \sum_{\{\sigma\}} \exp \left(J \sum_{\langle i,j \rangle} \sigma_i \sigma_j + H \sum_j \sigma_j \right)$$

- ▶ Analytic behaviour of $Z(J, H)$ w.r.t. complex variable H ?

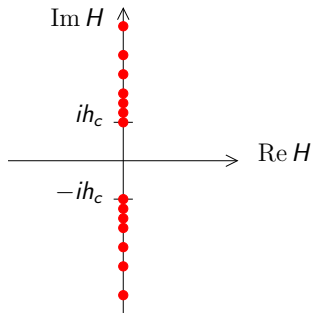
The Yang-Lee edge singularity

[Yang-Lee '52]

- ▶ Classical Ising model in magnetic field :

$$Z(J, H) = \sum_{\{\sigma\}} \exp \left(J \sum_{\langle i,j \rangle} \sigma_i \sigma_j + H \sum_j \sigma_j \right)$$

- ▶ Analytic behaviour of $Z(J, H)$ w.r.t. complex variable H ?
- ▶ Distribution of zeros of $Z(J, H)$ for $T > T_c$:



The Yang-Lee model as a minimal model of CFT

[Cardy '85]

- ▶ Scaling limit: fixed $T > T_c$ and $H = ih \rightarrow ih_c$ [Fisher '78]
 - ▶ Density of zeros $\rho(h, T) \sim (h - h_c)^\sigma$
 - ▶ Free energy $F(h, T) = \int dx \rho(x, T) \log(h - ix)$
 - ▶ Magnetisation $M(h, T) = -i\partial_h F(h, T) \sim (h - h_c)^\sigma$

The Yang-Lee model as a minimal model of CFT

[Cardy '85]

- ▶ Scaling limit: fixed $T > T_c$ and $H = ih \rightarrow ih_c$ [Fisher '78]
 - ▶ Density of zeros $\rho(h, T) \sim (h - h_c)^\sigma$
 - ▶ Free energy $F(h, T) = \int dx \rho(x, T) \log(h - ix)$
 - ▶ Magnetisation $M(h, T) = -i\partial_h F(h, T) \sim (h - h_c)^\sigma$
- ▶ Scaling features
 - ▶ No internal symm, only one non-trivial primary field ϕ
 - ▶ OPE: $\phi \times \phi \rightarrow \mathbf{1} + \phi$ Scaling law: $\sigma = \frac{h_\phi}{1-h_\phi}$
 - ▶ high-T exp. $\rightarrow \sigma \simeq -0.163 \rightarrow h_\phi \simeq -0.195$

The Yang-Lee model as a minimal model of CFT

[Cardy '85]

- ▶ Scaling limit: fixed $T > T_c$ and $H = ih \rightarrow ih_c$ [Fisher '78]
 - ▶ Density of zeros $\rho(h, T) \sim (h - h_c)^\sigma$
 - ▶ Free energy $F(h, T) = \int dx \rho(x, T) \log(h - ix)$
 - ▶ Magnetisation $M(h, T) = -i\partial_h F(h, T) \sim (h - h_c)^\sigma$
- ▶ Scaling features
 - ▶ No internal symm, only one non-trivial primary field ϕ
 - ▶ OPE: $\phi \times \phi \rightarrow \mathbf{1} + \phi$ Scaling law: $\sigma = \frac{h_\phi}{1-h_\phi}$
 - ▶ high-T exp. $\rightarrow \sigma \simeq -0.163 \rightarrow h_\phi \simeq -0.195$
- ▶ Corresponds to minimal model $\mathcal{M}(5, 2) =$

| | | | |
|----------|--------|--------|----------|
| 1 | ϕ | ϕ | 1 |
|----------|--------|--------|----------|

Central charge $c = -\frac{22}{5}$, dimensions $h_{\mathbf{1}} = 0$, $h_\phi = -\frac{1}{5}$

The Yang-Lee model as a minimal model of CFT

[Cardy '85]

- ▶ Scaling limit: fixed $T > T_c$ and $H = ih \rightarrow ih_c$ [Fisher '78]

- ▶ Density of zeros $\rho(h, T) \sim (h - h_c)^\sigma$

- ▶ Free energy $F(h, T) = \int dx \rho(x, T) \log(h - ix)$

- ▶ Magnetisation $M(h, T) = -i\partial_h F(h, T) \sim (h - h_c)^\sigma$

- ▶ Scaling features

- ▶ No internal symm, only one non-trivial primary field ϕ

- ▶ OPE: $\phi \times \phi \rightarrow \mathbf{1} + \phi$ Scaling law: $\sigma = \frac{h_\phi}{1-h_\phi}$

- ▶ high-T exp. $\rightarrow \sigma \simeq -0.163 \rightarrow h_\phi \simeq -0.195$

- ▶ Corresponds to minimal model $\mathcal{M}(5, 2) =$

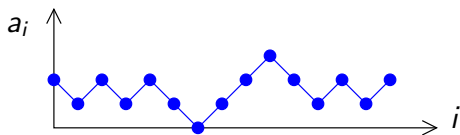
| | | | |
|--------------|--------|--------|--------------|
| $\mathbf{1}$ | ϕ | ϕ | $\mathbf{1}$ |
|--------------|--------|--------|--------------|

Central charge $c = -\frac{22}{5}$, dimensions $h_{\mathbf{1}} = 0$, $h_\phi = -\frac{1}{5}$

- ▶ Exactly solved RSOS model in same universality class

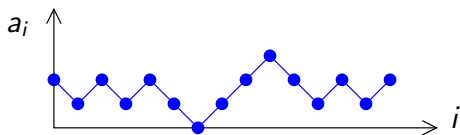
[Andrews-Baxter-Forrester '84]

The RSOS quantum chain



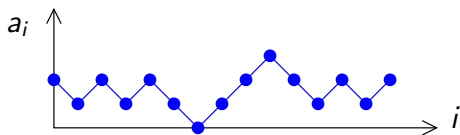
- Basis states: $|a_1, a_2, \dots, a_L\rangle$ with $\begin{cases} a_i \in \{1, \dots, p-1\} \\ |a_i - a_{i+1}| = 1 \end{cases}$

The RSOS quantum chain



- ▶ Basis states: $|a_1, a_2, \dots, a_L\rangle$ with $\begin{cases} a_i \in \{1, \dots, p-1\} \\ |a_i - a_{i+1}| = 1 \end{cases}$
- ▶ Hamiltonian: $H = -\sum_{i=1}^L e_i$ with PBC

The RSOS quantum chain



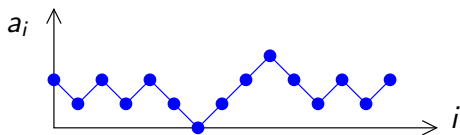
- ▶ Basis states: $|a_1, a_2, \dots, a_L\rangle$ with $\begin{cases} a_i \in \{1, \dots, p-1\} \\ |a_i - a_{i+1}| = 1 \end{cases}$

- ▶ Hamiltonian: $H = -\sum_{i=1}^L e_i$ with PBC

- ▶ Action on states:

$$e_i |\dots a_{i-1}, a_i, a_{i+1} \dots\rangle = \delta_{a_{i-1}, a_{i+1}} \sum_{a'_i} \frac{\sin \lambda a'_i}{\sin \lambda a_i} |\dots a_{i-1}, a'_i, a_{i+1} \dots\rangle$$

The RSOS quantum chain



- ▶ Basis states: $|a_1, a_2, \dots, a_L\rangle$ with $\begin{cases} a_i \in \{1, \dots, p-1\} \\ |a_i - a_{i+1}| = 1 \end{cases}$

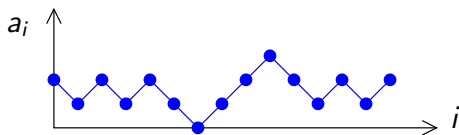
- ▶ Hamiltonian: $H = -\sum_{i=1}^L e_i$ with PBC

- ▶ Action on states:

$$e_i |\dots a_{i-1}, a_i, a_{i+1} \dots\rangle = \delta_{a_{i-1}, a_{i+1}} \sum_{a'_i} \frac{\sin \lambda a'_i}{\sin \lambda a_i} |\dots a_{i-1}, a'_i, a_{i+1} \dots\rangle$$

- ▶ $\lambda = \frac{\pi(p-q)}{p}$ with $p > q$ coprime \rightarrow scaling limit = $\mathcal{M}(p, q)$

The RSOS quantum chain



- ▶ Basis states: $|a_1, a_2, \dots, a_L\rangle$ with $\begin{cases} a_i \in \{1, \dots, p-1\} \\ |a_i - a_{i+1}| = 1 \end{cases}$

- ▶ Hamiltonian: $H = -\sum_{i=1}^L e_i$ with PBC

- ▶ Action on states:

$$e_i |\dots a_{i-1}, a_i, a_{i+1} \dots\rangle = \delta_{a_{i-1}, a_{i+1}} \sum_{a'_i} \frac{\sin \lambda a'_i}{\sin \lambda a_i} |\dots a_{i-1}, a'_i, a_{i+1} \dots\rangle$$

- ▶ $\lambda = \frac{\pi(p-q)}{p}$ with $p > q$ coprime \rightarrow scaling limit = $\mathcal{M}(p, q)$
- ▶ The YL case : $p = 5, \quad \lambda = \frac{3\pi}{5}$

EEs of a non-unitary model

see also [Bianchini,Castro-Alvaredo,Doyon,Levi,Ravanini '14]

- ▶ $h_1 = 0$: the conformally invariant state is $|\mathbf{1}\rangle$.

EEs of a non-unitary model

see also [Bianchini,Castro-Alvaredo,Doyon,Levi,Ravanini '14]

- ▶ $h_{\mathbf{1}} = 0$: the conformally invariant state is $|\mathbf{1}\rangle$.
- ▶ $h_{\phi} = -1/5 < 0$: the ground state is $|\phi\rangle$.

EEs of a non-unitary model

see also [Bianchini,Castro-Alvaredo,Doyon,Levi,Ravanini '14]

- ▶ $h_{\mathbf{1}} = 0$: the conformally invariant state is $|\mathbf{1}\rangle$.
- ▶ $h_{\phi} = -1/5 < 0$: the ground state is $|\phi\rangle$.
- ▶ Dimension of twisted operator τ_{ϕ} : $\hat{h}_{\phi} = \frac{c}{24}(N - 1/N) + \frac{h_{\phi}}{N}$

EEs of a non-unitary model

see also [Bianchini,Castro-Alvaredo,Doyon,Levi,Ravanini '14]

- ▶ $h_{\mathbf{1}} = 0$: the conformally invariant state is $|\mathbf{1}\rangle$.
- ▶ $h_{\phi} = -1/5 < 0$: the ground state is $|\phi\rangle$.
- ▶ Dimension of twisted operator τ_{ϕ} : $\hat{h}_{\phi} = \frac{c}{24}(N - 1/N) + \frac{h_{\phi}}{N}$
- ▶ τ_{ϕ} is more relevant than $\tau_{\mathbf{1}}$

EEs of a non-unitary model

see also [Bianchini,Castro-Alvaredo,Doyon,Levi,Ravanini '14]

- ▶ $h_{\mathbf{1}} = 0$: the conformally invariant state is $|\mathbf{1}\rangle$.
- ▶ $h_{\phi} = -1/5 < 0$: the ground state is $|\phi\rangle$.
- ▶ Dimension of twisted operator τ_{ϕ} : $\hat{h}_{\phi} = \frac{c}{24}(N - 1/N) + \frac{h_{\phi}}{N}$
- ▶ τ_{ϕ} is more relevant than $\tau_{\mathbf{1}}$
- ▶ Consider $T = 0$, one-interval EE with PBC.
⇒ Conformal map: $x = \exp(2i\pi\ell/L)$

EEs of a non-unitary model

see also [Bianchini,Castro-Alvaredo,Doyon,Levi,Ravanini '14]

- ▶ $h_{\mathbf{1}} = 0$: the conformally invariant state is $|\mathbf{1}\rangle$.
- ▶ $h_{\phi} = -1/5 < 0$: the ground state is $|\phi\rangle$.
- ▶ Dimension of twisted operator τ_{ϕ} : $\hat{h}_{\phi} = \frac{c}{24}(N - 1/N) + \frac{h_{\phi}}{N}$
- ▶ τ_{ϕ} is more relevant than $\tau_{\mathbf{1}}$
- ▶ Consider $T = 0$, one-interval EE with PBC.
⇒ Conformal map: $x = \exp(2i\pi\ell/L)$
- ▶ EE in state $|\mathbf{1}\rangle$:

$$\langle \tau_{\phi}(\mathbf{1}) \tilde{\tau}_{\phi}(x, \bar{x}) \rangle = |x - 1|^{-4\hat{h}_{\phi}} = \left(2 \sin \frac{\pi\ell}{L} \right)^{-4\hat{h}_{\phi}}$$

EEs of a non-unitary model

see also [Bianchini,Castro-Alvaredo,Doyon,Levi,Ravanini '14]

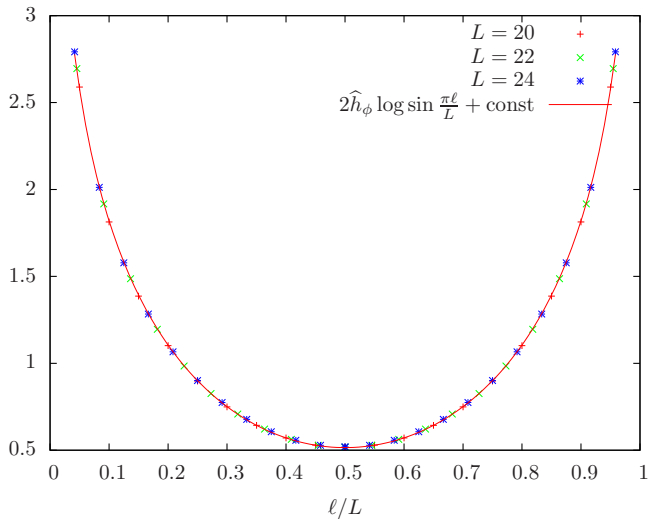
- ▶ $h_{\mathbf{1}} = 0$: the conformally invariant state is $|\mathbf{1}\rangle$.
- ▶ $h_{\phi} = -1/5 < 0$: the ground state is $|\phi\rangle$.
- ▶ Dimension of twisted operator τ_{ϕ} : $\hat{h}_{\phi} = \frac{c}{24}(N - 1/N) + \frac{h_{\phi}}{N}$
- ▶ τ_{ϕ} is more relevant than $\tau_{\mathbf{1}}$
- ▶ Consider $T = 0$, one-interval EE with PBC.
⇒ Conformal map: $x = \exp(2i\pi\ell/L)$
- ▶ EE in state $|\mathbf{1}\rangle$:

$$\langle \tau_{\phi}(\mathbf{1}) \tilde{\tau}_{\phi}(x, \bar{x}) \rangle = |x - 1|^{-4\hat{h}_{\phi}} = \left(2 \sin \frac{\pi\ell}{L} \right)^{-4\hat{h}_{\phi}}$$

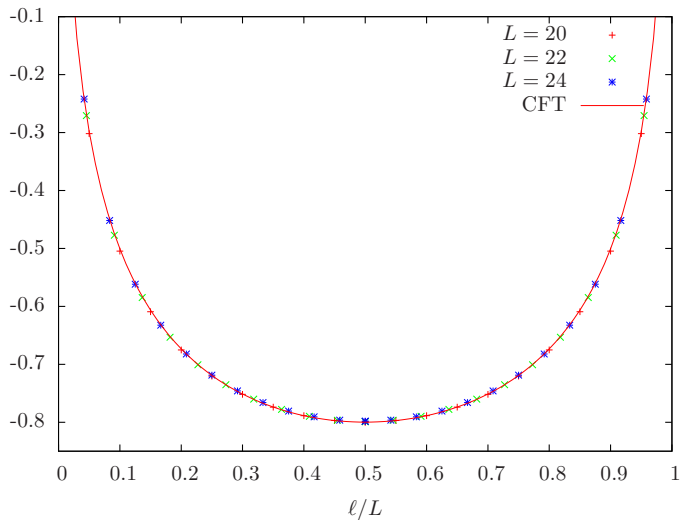
- ▶ EE in ground state $|\phi\rangle$:

$$\langle \Phi(\infty) \tau_{\phi}(\mathbf{1}) \tilde{\tau}_{\phi}(x, \bar{x}) \Phi(0) \rangle = G(x, \bar{x}) \quad \text{with } \Phi = \phi^{\otimes N}$$

$N = 3$ Rényi entropy in state $|1\rangle$



$N = 3$ Rényi entropy in the ground state $|\phi\rangle$



Conclusions and perspectives

► Current results

- “Standard” computations of EE: only $g = 0$ **or** $c \in \{\frac{1}{2}, 1\}$
- New approach based on \mathbb{Z}_N -orbifold of Virasoro algebra
- Works for minimal models (twist has two null vectors)
- Applied at $g = 0$ for YL : gives non-trivial $\langle \dots \tau_\phi \tilde{\tau}_\phi \dots \rangle$
- Tested at $g = 1$ for YL+Ising : recover $\{\chi_j\}$

► Future work

- Find systematic derivation of differential equations ?
- Understand fusion rules in \mathbb{Z}_N -orbifold ?
- Construct Coulomb-Gas formalism for conformal blocks in \mathbb{Z}_N -orbifold ? [joint with O. Blondeau-Fournier (Laval)]

Thank you!